

Energy Difference between the Classical and the Nonclassical 2-Norbornyl Cation in Solution. A Combined *ab Initio*–Monte Carlo Aqueous Solution Study

Peter R. Schreiner,^{1a,c} Daniel L. Severance,^{1b,†}
William L. Jorgensen,^{1b} Paul von Ragué Schleyer,^{*,1a,c} and
Henry F. Schaefer III^{1a}

Center for Computational Quantum Chemistry
University of Georgia, Athens, Georgia 30602
Department of Chemistry, Yale University
New Haven, Connecticut 06520-8107
Computer Chemistry Center, Institut für Organische Chemie
der Universität Erlangen-Nürnberg
Henkestr. 42, D-91054 Erlangen, Germany

Received September 6, 1994

While theory and experiment agree on the nonclassical, symmetrically bridged minimum energy structure (C_s , **1**) for the 2-norbornyl cation,² it is not always straightforward to compare experimental data obtained mostly in solution with computed gas phase values. The classical 2-norbornyl cation (**2**) is a transition structure² which can only be structurally characterized by computation. To close the gap between theory and experiment with respect to the 2-norbornyl cation problem, we report our study on the effects of solvation on the classical (**2**) and nonclassical ion (**1**) by employing *ab initio* MO methods and statistical Monte Carlo (MC) solution simulations.

Cation **1** is more stable than **2** by 13.6 kcal mol⁻¹ [MP4-(fc)SDQ/6-31G**/MP2(full)/6-31G* + ZPVE], but the difference in the transition state energies for the 2-*exo* and 2-*endo* norbornyl solvolyses³ is reduced to less than half of this value (4–7 kcal mol⁻¹ from experiment,⁴ 6 kcal mol⁻¹ from theory).⁵ We will show in another place⁵ that bridging does lag behind ionization,⁶ and that this effect should be mainly responsible for the aforementioned differences in activation energies. The hypothesis that the relative energies of anchimerically assisted solvolysis reactions are very similar in the gas phase and in solution, implying that relative carbocation stabilities are the same in both media, is probed in the present work by including solvent effects in the energy computations.

We chose water as our solvent since we expect the difference in solvation energies to be large in highly polar media. To obtain the differences in free energies of hydration ($\Delta\Delta G_{\text{hyd}}$, Figure 1), we perturbed⁷ between the MP2(full)/6-31G* optimized⁸ cation structures **1** and **2** using statistical perturbation theory⁹ as implemented in the BOSS program.¹⁰ The intermolecular interactions are represented by Coulomb and Lennard-Jones terms with all atoms explicit.¹¹ For water, the TIP4P model was adopted,¹² while standard OPLS Lennard-Jones parameters were used for the solute.¹¹ Charges were obtained from MP2(full)/6-31G* – CHELPG¹³ calculations. The simu-

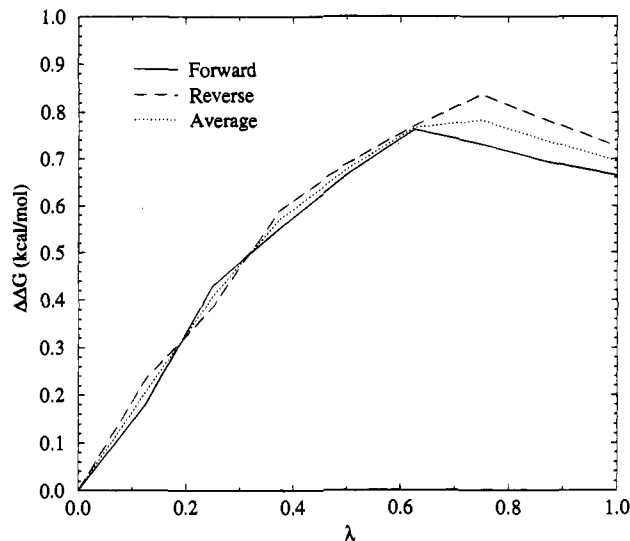
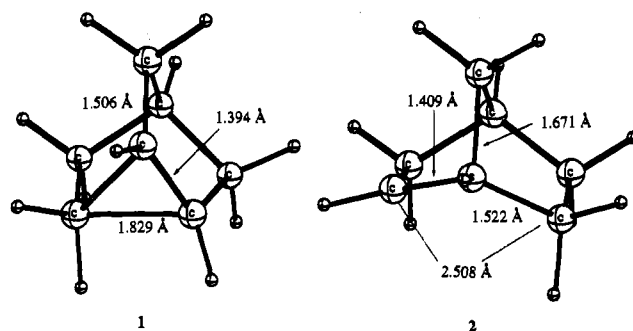


Figure 1. Free energy of solvation ($\Delta\Delta G$) vs coupling parameter λ ($\lambda = 0$ corresponds to **2**, $\lambda = 1$ to **1**).⁷



lations were performed in the NPT ensemble at 25 °C and 1 atm with Metropolis and preferential sampling.⁹ The system consisted of 506 water molecules and the solute in a periodic cell, ca. 25 Å on a side. Each of the six simulations for the mutation entailed 10⁶ configurations of equilibration and 2 × 10⁶ configurations of averaging. Solute–water interactions were included for water molecules with an oxygen within 10 Å of any solute atom, and the water–water cutoff was set at 10 Å on the basis of the O–O distance.

First, we established a reference point between the *ab initio* and BOSS force field results by optimizing the structures and energies of the nonclassical ion–water complex.¹⁴ Both methods yield remarkably similar structures (**3** and **4**, both C_s)

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[†] Current address: Syntex Discovery Research RGW-002, 3401 Hillview Ave, Palo Alto, CA 94303.

(1) (a) University of Georgia. (b) Yale University. (c) Universität Erlangen-Nürnberg.

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(3) Winstein, S.; Trifan, D. S. *J. Am. Chem. Soc.* **1949**, *71*, 2953. Winstein, S.; Trifan, D. S. *J. Am. Chem. Soc.* **1952**, *74*, 1147, 1154.

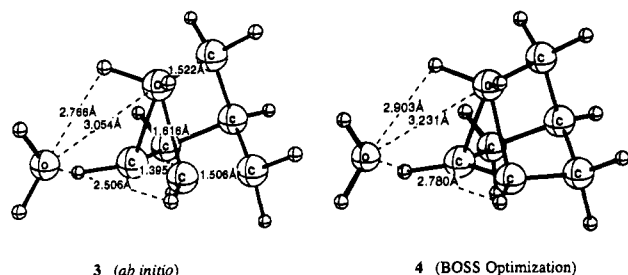
(4) For the most recent review on the norbornyl solvolysis, see: Lenoir, D.; Apeloig, Y.; Arad, D.; Schleyer, P. v. R. *J. Org. Chem.* **1988**, *53*, 661 and references cited therein.

(5) At MP2(fc)/6-31G**//HF/6-31G* + ZPVE: Schreiner, P. R.; Schleyer, P. v. R.; Schaefer, H. F. Submitted for publication.

(6) Winstein, S. *J. Am. Chem. Soc.* **1965**, *87*, 381.

(7) A linear interpolation of the internal coordinates for the conversion of **1** into **2** was performed as a function of the coupling parameter λ .

and complexation energies¹⁵ [MP4(fc)SDQ/6-31+G*/MP2-(full)/6-31G* = 12.3 kcal mol⁻¹; BOSS = 10.4 kcal mol⁻¹], although the computational approaches are entirely different.

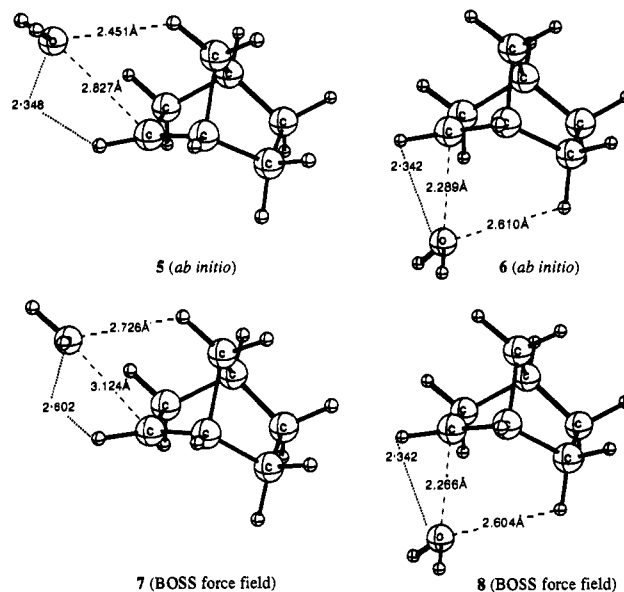


The $\Delta\Delta G_{\text{hyd}}$ (Figure 1) profile summarizes the present MC solution simulations. To our surprise, the classical cation (**2**) is *not* considerably more stabilized by hydration than the nonclassical cation (**1**), as the $\Delta\Delta G_{\text{hyd}}$ is only 0.7 kcal mol⁻¹ in favor of **2**. Partial (allowing only the water molecule to be optimized while keeping the cation moiety fixed) *ab initio* (**5** and **6**) and BOSS force field optimizations (**7** and **8**) of the possible complexes of **2** with a water molecule yield a similar result.

While the complexation energy of **3** is 12.3 kcal mol⁻¹, we find an interaction energy of 12.8 kcal mol⁻¹ for **5** and 8.5 kcal mol⁻¹ for **6** at MP4(fc)SDQ/6-31+G*/MP2(full)/6-31G*. Thus, the preferentially more favored *exo* complex of **2** with water is only 0.5 kcal mol⁻¹ more stabilized than the water complex of **1**. One might have speculated that **2** has a more localized positive charge and would therefore benefit more from solvation than the delocalized **1**. This assumption does not hold as the MP2(full)/6-31G* - CHELPG¹¹ charges show very similar partial charges (+0.10 to +0.15) on the hydrogens close to the cationic centers. Moreover, the sums of all hydrogen charges are comparable (+1.2 for **2**; +1.1 for **1**). The sp² carbon in **2** does not have a full positive charge (+0.36). Since the water-cation binding appears to feature interactions with the cation's hydrogens, the energies are quite similar.

(14) We used the MP2(full)/6-31G* optimized geometry for the MC simulation, i.e., only the position and energy of the water molecule were optimized. One may expect geometrical changes upon relaxation of the cation geometries in the solvent which would probably be more important for the less stable classical ion. However, *ab initio* optimizations within the MC simulation are currently not feasible, and the relative stabilities are not expected to change much upon reoptimization in solution.

(15) In a more rigorous treatment, one should include entropic effects and ZPVE corrections in the *ab initio* complexation energy, but frequency calculations were computationally not feasible for **3**. Although we have used diffuse functions for our energy evaluations to overcome (see: Feller, *D. J. Chem. Phys.* **1992**, *96*, 6104) most of the basis set superposition error (BSSE), the stability of **3** may still be somewhat overestimated. In short, 12.3 kcal mol⁻¹ seems to be an upper limit for the complexation energy.



In conclusion, the classical 2-norbornyl cation (**2**) is not significantly more stabilized in aqueous solution than the nonclassical ion (**1**). Thus both the gas phase *ab initio* computations and Monte Carlo solution simulation come to the same result. In less polar solvents, the solvation energies are expected to be even more similar. Our results confirm that the nonclassical form **1** of the 2-norbornyl cation is the only stable form in the gas phase and in solution. The classical form **2** is unlikely to be involved in solvolysis reactions. As a consequence, the differences in rates of solvolysis for 2-*exo* and 2-*endo* norbornyl derivatives can only be explained in terms of the differences in the *exo* vs *endo* transition states.⁵

Acknowledgment. This research was supported by the Fonds der Chemischen Industrie (fellowship for the first author), the Convex Computer Corporation, the U.S. Department of Energy (Office of Basic Energy Sciences, Division of Chemical Sciences, Fundamental Interactions Branch, Grant DE-FG05-94ER14428), and the National Science Foundation.

Supplementary Material Available: Cartesian coordinates, charges, and values for ϵ for the Lennard-Jones potentials for **1** and **2** (4 pages). This material is contained in many libraries on microfiche, immediately follows this article in the microfilm version of the journal, and can be ordered from the ACS; see any current masthead page for ordering information.

JA942946C